

PHYS 5770 Grav Theory Spring 2008. Problem Set 5. Due Tue 15 Apr

**1. Tidal forces falling into a Schwarzschild black hole**

In the Gullstrand-Painlevé tetrad, the non-zero components of the tetrad-frame Riemann tensor are

$$-\frac{1}{2}R_{trtr} = R_{t\theta t\theta} = R_{t\phi t\phi} = -R_{r\theta r\theta} = -R_{r\phi r\phi} = \frac{1}{2}R_{\theta\phi\theta\phi} = C \quad (1.1)$$

where

$$C = -\frac{M}{r^3} \quad (1.2)$$

is the Weyl scalar.

**(a) Tidal forces**

By construction of the Gullstrand-Painlevé tetrad, a person who falls radially from zero velocity at infinity is at rest in the Gullstrand-Painlevé tetrad, with tetrad-frame 4-velocity  $u^m = \{1, 0, 0, 0\}$ . From the equation of geodesic deviation

$$\frac{D^2\delta\xi_m}{D\tau^2} + R_{klmn}\delta\xi^k u^l u^n = 0 \quad (1.3)$$

deduce the tidal acceleration on the person in the radial and transverse directions. Does the tidal acceleration stretch or compress? [Hint: The equation of geodesic deviation gives the proper acceleration between two points a small distance  $\delta\xi^m$  apart, where  $\xi^m$  are the locally inertial coordinates of the tetrad frame. Notice that this problem is much easier to solve with tetrads than with the traditional coordinate approach.]

**(b) Choose a black hole to fall into**

What is the mass of the black hole for which the tidal acceleration  $M/r^3$  is 1 gee per meter at the horizon? If you wanted to fall through the horizon of a black hole without first being torn apart, what mass of black hole would you choose? [Hint: 1 gee is the gravitational acceleration at the surface of the Earth.]

**(c) Time to die**

In problem set 3 you showed that the proper time to free-fall radially from radius  $r$  to the singularity of a Schwarzschild black hole is

$$\tau = \frac{\sqrt{2}}{3} \sqrt{\frac{r^3}{M}}. \quad (1.4)$$

How long, in seconds, does it take to fall to the singularity from the place where the tidal acceleration is 1 gee per meter?

## 2. Constant density star

A variation of the following calculation (done in a very different way) of a static spherically symmetric gravitating system convinced Einstein (1939, Ann. Math. 40, 922) that black holes cannot exist.

In a spherically symmetric static spacetime, Einstein's equations reduce to an equation for the mass  $M$  interior to  $r$

$$\frac{dM}{dr} = 4\pi r^2 \rho, \quad (2.1)$$

and to the Volkov-Oppenheimer equation of hydrostatic equilibrium

$$\frac{dp}{dr} = -\frac{(\rho + p)(M + 4\pi r^3 p)}{r^2(1 - 2M/r)}. \quad (2.2)$$

### (a) Interior mass

Suppose that the density  $\rho$  is constant. From equation (2.1) obtain an expression for the interior mass  $M$  as a function of radius  $r$  and the density  $\rho$ . [Hint: This is easy.]

### (b) Hydrostatic equilibrium

Given your expression for  $M$ , show that the Volkov-Oppenheimer equation rearranges to

$$\int_{p_c} \frac{dp}{(\rho + p)(\rho + 3p)} = - \int_0^r \frac{4\pi r dr}{3 - 8\pi r^2 \rho} \quad (2.3)$$

where  $p_c$  is the central pressure, where the radius is zero,  $r = 0$ .

### (c) Solve

Integrate equation (2.3). From the integral evaluated at the edge of the star, where the pressure is zero,  $p = 0$ , and the radius is the stellar radius,  $r = R_*$ , argue that

$$\frac{\rho + 3p_c}{\rho + p_c} = \sqrt{\frac{1}{1 - 2M_*/R_*}} \quad (2.4)$$

where  $M_* \equiv \frac{4}{3}\pi\rho R_*^3$  is the total mass of the star.

### (d) Limits

From the condition that the central pressure be positive and finite,  $0 < p_c < \infty$ , deduce that

$$0 < \frac{2M_*}{R_*} < \frac{8}{9}. \quad (2.5)$$

### (e) Comment

Comment on what equation (2.5) implies physically. [Hint: What is the Schwarzschild radius?]